

## LAGRANGIAN SINGLE-PARTICLE STATISTICS OF FLUID TURBULENCE

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**Abstract** Based on observations of experimental and numerical data and theoretical considerations, we question the dimensional scaling of the second-order Lagrangian velocity structure function  $\langle [v(t+\tau) - v(t)]^2 \rangle \sim \epsilon\tau$ . We show that that state-of-the-art Lagrangian data up to  $R_\lambda = O(10^3)$  are consistent with a scaling relation  $\langle [v(t+\tau) - v(t)]^2 \rangle \sim \tau^{1-\mu}$  with the small correction on exponent  $\mu \approx 0.1$ . We also discuss further implications of this breakdown of the dimensional scaling relation.

For three-dimensional fluid turbulence, kinetic energy is supplied at large scales and then transferred to smaller scales until it is dissipated by viscosity. The well-known Kármán-Howarth-Kolmogorov equation establishes an exact relation between the rate of energy transfer (energy cascade),  $\epsilon$ , and the third-order Eulerian velocity structure function in space [1, 2]:

$$\langle \{[\mathbf{u}(\mathbf{x} + \mathbf{r}) - \mathbf{u}(\mathbf{x})] \cdot (\mathbf{r}/r)\}^3 \rangle = -\frac{4}{5}\epsilon r, \quad (1)$$

where  $r = |\mathbf{r}|$  is the distance between the two points separated by vector  $r$  at which the velocities  $\mathbf{u}(\mathbf{x} + \mathbf{r})$  and  $\mathbf{u}(\mathbf{x})$  are measured simultaneously. This exact “4/5-law” has been the foundations of nearly all the subsequent theoretical and phenomenological work on Eulerian statistics of turbulence.

Lagrangian properties of turbulence, *i.e.*, how turbulence looks like when travelling with fluid particles in the flow, have attracted increasing interest in the last two decades, due to a combination of theoretical breakthroughs [3, 4] and the rapid advances in experimental [5, 6, 7, 8] and numerical techniques [9, 10]. However, unlike in the cases of Eulerian statistics, there is no known exact result on inertial range Lagrangian single-particle statistics. It is commonly assumed that in the inertial range the dimensional scaling relation of the second-order Lagrangian velocity structure function holds [1, 11]:

$$D_2(\tau) \equiv \langle \delta_\tau^2 v \rangle \equiv \langle [v(t+\tau) - v(t)]^2 \rangle = C_0 \epsilon \tau, \quad (\tau_\eta \ll \tau \ll T_L) \quad (2)$$

where  $v(t)$  is the velocity following a fluid particle in turbulence,  $C_0$  is assumed to be a universal dimensionless constant,  $\tau_\eta$  and  $T_L$  are the Kolmogorov and integral time scales of the flow, respectively. The reason that the second-order structure function  $D_2(t)$  is treated specially is that there is no apparent intermittency correction on the scaling of  $\tau$  in this dimensional argument. Considerable effort has been devoted to verify Eq. (2) or to identify the numerical values of  $C_0$  [9]. However, no convincing inertial scaling relation as predicted by Eq. (2) has been observed in state-of-the-art experimental and numerical simulation data [10, 12].

Based on theoretical considerations and observations of experimental and numerical data, we argue that the dimensional scaling Eq. (2) might be fundamentally flawed [13]. For example, if Eq. (2) holds, then the following will be true:

$$\frac{dD_2}{d\tau} = 2\langle a\delta_\tau v \rangle = C_0 \epsilon \quad (3)$$

in the inertial range. However, as shown in Figure 1,  $dD_2/d\tau$  observed from experiments and numerical simulations reaches maximum at approximately  $2\tau_\eta$  and then decreases nearly exponentially, without appreciable plateau range. Further investigation of the acceleration spectra obtained from numerical simulations indicate that a small correction on the scaling exponent given in Eq. (2) better fits the data:  $D_2 \sim \tau^{1-\mu}$  with  $\mu \approx 0.1$ .

This small correction, if it holds true at larger Reynolds numbers, implies that the choice of available parameters in the dimensional argument leading to Eq. (2) is incorrect. If  $\epsilon$  is not the right parameter, then what is the alternative? A deeply related question is: How are Lagrangian single-particle statistics related to  $\epsilon$ , the energy cascade through *spatial scales*? We will discuss these in the presentation.

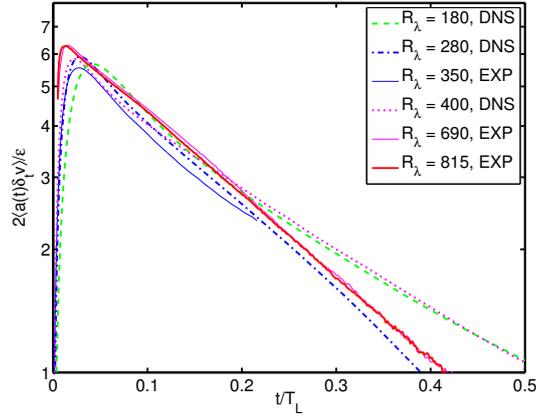
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**Figure 1.** Derivative of the second order velocity structure function,  $\frac{dD_2}{d\tau} = 2\langle a\delta_\tau v \rangle$  plotted versus  $\tau/T_L$ . Eq. (2) predicts a plateau in the inertial range.

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